

Tidal forces and mode resonances in compact binaries

Sayan Chakrabarti¹ T rence Delsate² Jan Steinhoff³



¹Indian Institute of Technology
Guwahati, India



²UMons
Mons, Belgium



T CNICO
LISBOA

³Instituto Superior T cnico
Lisbon, Portugal

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through STE 2017 “Resonances of quasinormal modes and orbital motion in
general relativistic compact binaries”

Important compact object: Neutron star

`pics/neutronstar.pdf`

Neutron star picture by D. Page

www.astroscu.unam.mx/neutrones/

„Lab“ for various areas in physics

- magnetic field, plasma
- **crust** (solid state)
- superfluidity
- superconductivity
- unknown matter in core
(condensate of quarks, hyperons,
kaons, pions, ... ?)

other important objects:

- magnetars, quark stars
- black holes, boson stars

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Tidal forces in binaries and resonances

tidal forces in inspiraling binaries \longleftrightarrow oscillation modes of neutron stars
 \Rightarrow resonances!

resonances probably
relevant for short
gamma-ray bursts

[Tsang et.al., PRL 108 (2012) 011102]

Swift/BAT, nasa.gov

`pics/double_pulsar.pdf`

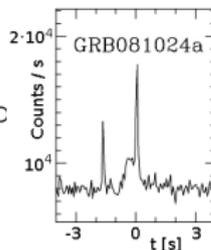
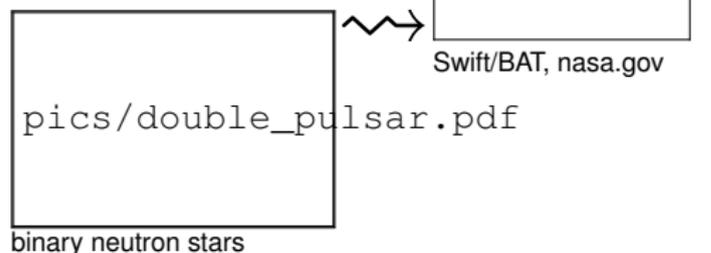
binary neutron stars

mode spectrum from
gravitational waves:
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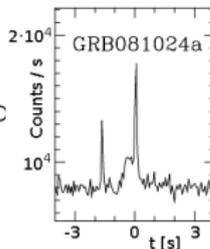
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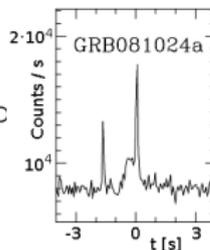
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Einstein Telescope

pics/spectral.p

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Tidal forces in Newtonian gravity

Simple case: linear perturbation of nonrotating barotropic stars

temperature-independent equation of state

see e.g. [Press, Teukolsky, ApJ **213** (1977) 183]

Displacement $\vec{\xi} :=$ perturbed – unperturbed location of fluid elements

$$\ddot{\vec{\xi}} + \mathcal{D}\vec{\xi} = (\text{external forces})$$

$$\mathcal{D}\vec{\xi} := -\vec{\nabla} \left\{ \left[\frac{c_s^2}{\rho_0} + 4\pi G\Delta^{-1} \right] \vec{\nabla} \cdot (\rho_0 \vec{\xi}) \right\}$$

ρ_0 : unperturbed mass density

c_s : speed of sound

G : Newton constant

Properties of operator \mathcal{D} :

- contains differential operators $\vec{\nabla}$
- and also integral operator Δ^{-1}
- linear, nonlocal
- spherical symmetric
- Hermitian w.r.t. compact measure $dm_0 := \rho_0 d^3x$

[Chandrasekhar, ApJ **139** (1964) 664–674]

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Normal modes (NM) in Newtonian gravity

Eigenfunctions of \mathcal{D} are the **normal oscillation modes** of the star

\mathcal{D} is Hermitian, compact integration measure $dm_0 := \rho_0 d^3x$

⇒ **Discrete, real eigenvalues or oscillation frequencies** $\omega_{n\ell}$

Eigenvalue equation:

$$\mathcal{D}\vec{\xi}_{n\ell m}^{\text{NM}} = \omega_{n\ell}^2 \vec{\xi}_{n\ell m}^{\text{NM}}$$

Orthonormalization:

$$\int dm_0 \vec{\xi}_{n'\ell'm'}^{\text{NM}\dagger} \vec{\xi}_{n\ell m}^{\text{NM}} = \delta_{n'n} \delta_{\ell'\ell} \delta_{m'm}$$

Indices $\{\ell, m\}$ from spherical harmonics

Decomposition in terms of **mode amplitudes** $A_{n\ell m}(t)$

$$\vec{\xi} = \sum_{n\ell m} A_{n\ell m}(t) \vec{\xi}_{n\ell m}^{\text{NM}}(\vec{x})$$

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Amplitude formulation

Displacement field $\vec{\xi}$ specified by

$$\begin{aligned}\ddot{\vec{\xi}} + \mathcal{D}\vec{\xi} &= (\text{external forces}) \\ &= -\vec{\nabla}\Phi_{\text{ext}}\end{aligned}$$

Insert mode decomposition

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and project onto orthonormal basis $\vec{\xi}_{nlm}^{\text{NM}}$

Result: uncoupled forced **harmonic oscillators**

$$\begin{aligned}\ddot{A}_{nlm} + \omega_{nl}^2 A_{nlm} &= f_{nlm} \\ f_{nlm} &:= - \int dm_0 \vec{\xi}_{nlm}^{\text{NM}} \cdot \vec{\nabla}\Phi_{\text{ext}}\end{aligned}$$

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Radial/angular split of integral (spherical harmonics)

$I_{nl} \sim$ radial integration part of f_{nlm}
overlap integral

$f_{nlm} \propto I_{nl} \times (\ell\text{-pole of } \Phi_{\text{ext}})$

Overlap integrals I_{nl}

- Coupling constants between modes to external field
- Determine energy exchange between orbital motion and modes
- Together with frequencies ω_{nl} they define the **gravito-spectrum** (measurable through gravitational waves)

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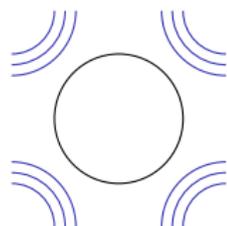
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Different perspective on the Newtonian theory

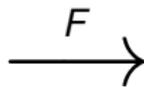
Chakrabarti, Delsate, Steinhoff, arXiv:1306.5820

Example: quadrupolar sector $\ell = 2$

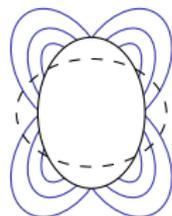


external quadrupolar field

linear response



→ deformation →



quadrupolar response

quadrupolar
response:

poles \Rightarrow resonances!

$$F = \sum_n \frac{I_{n\ell}^2}{\omega_{n\ell}^2 - \omega^2}$$

$\omega_{n\ell}$: mode frequency

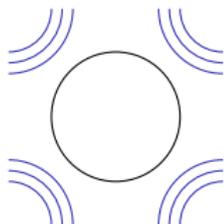
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R : radius

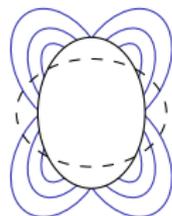
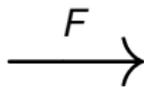
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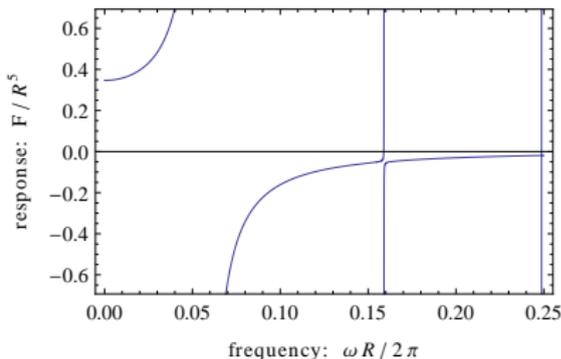
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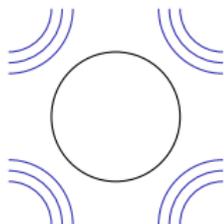
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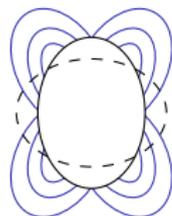
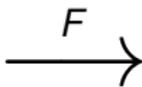
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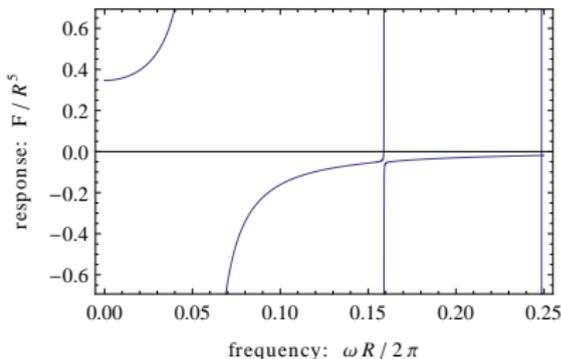
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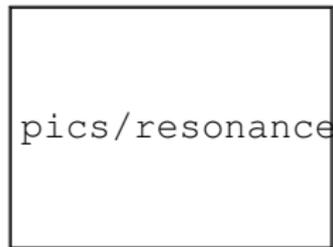
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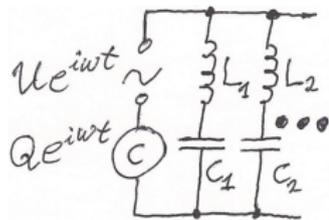
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Tacoma Bridge

Analogy with electronics

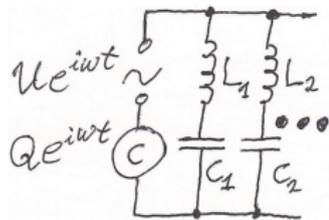


$$\begin{aligned}\frac{Q}{U} &= \frac{1}{i\omega Z} \\ &= \sum_n \frac{\frac{1}{L_n}}{\frac{1}{C_n L_n} - \omega^2}\end{aligned}$$

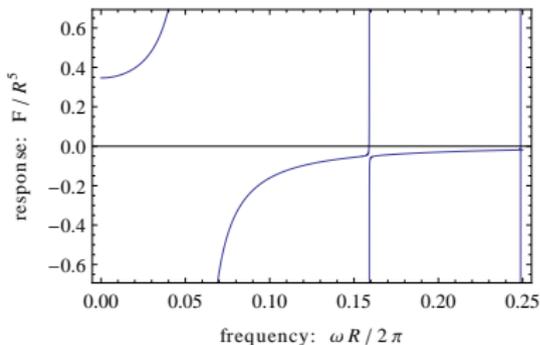
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E: external tidal field

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Philosophy of the approach

Starting point: single object, e.g., neutron star

Idea

Multipoles describe compact object on macroscopic scale

\longleftrightarrow

| | | |
|-------------------------------|-----------------------|--------------------------|
| state variables (p, V, T) | \longleftrightarrow | multipoles (M, S, Q) |
| equations of state | \longleftrightarrow | effective action |
| correlation | \longleftrightarrow | response |

Approximations for binary system using effective theory:

- Effective point-particle action with dynamical multipoles
- Response functions (propagators) for multipoles

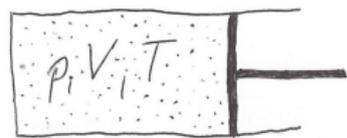
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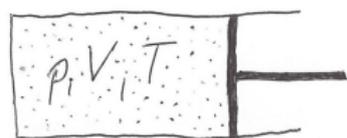
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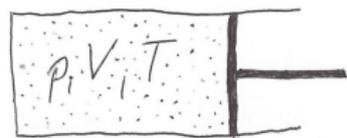
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Newtonian theory in terms of effective action and response functions:

- Alternative computation of overlap integral through fit of F
- Immediately generalizes to more complicated situations
- Can be generalized to the relativistic case

Relativistic tidal interactions

Status: expansion around adiabatic case

$$F(\omega) = 2\mu_2 + i\omega\lambda + 2\omega^2\mu'_2 + \mathcal{O}(\omega^3)$$

- μ_2 : 2nd kind relativistic Love number [Hinderer, ApJ **677** (2008) 1216; Damour, Nagar, PRD **80** (2009) 084035; Binnington, Poisson, PRD **80** (2009) 084018]
- λ : absorption [Goldberger, Rothstein, PRD **73** (2006) 104030]
- μ'_2 : beyond adiabatic, **not yet computed** [Bini, Damour, Faye, PRD **85** (2012) 124034]

Motivation:

- Adiabatic tidal effects may not be sufficient [Maselli, Gualtieri, Pannarale, Ferrari, PRD **86** (2012) 044032]
- Definition of **relativistic overlap integrals**
- Resonances between oscillation modes and orbital motion:
 - Numerical simulations of binary neutron stars for eccentric orbits [Gold, Bernuzzi, Thierfelder, Brüggmann, Pretorius, PRD **86** (2012) 121501]
 - **Shattering of neutron star crust** [Tsang, Read, Hinderer, Piro, Bondarescu, PRL **108** (2012) 011102]
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Relativistic case: generic problems

- Analog to Hermitian operator \mathcal{D} not available in the relativistic case
- Substitute neutron star by point particle reproducing large-scale field



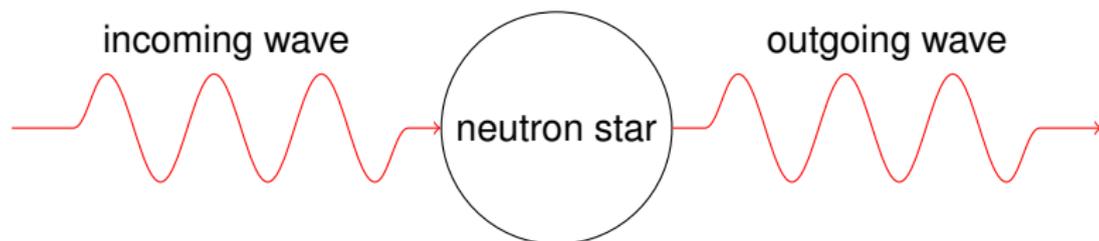
- Need to handle an inhomogeneous Regge-Wheeler equation with **effective point-particle source S** representing a neutron star

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- How to construct solutions corresponding to external field and response?

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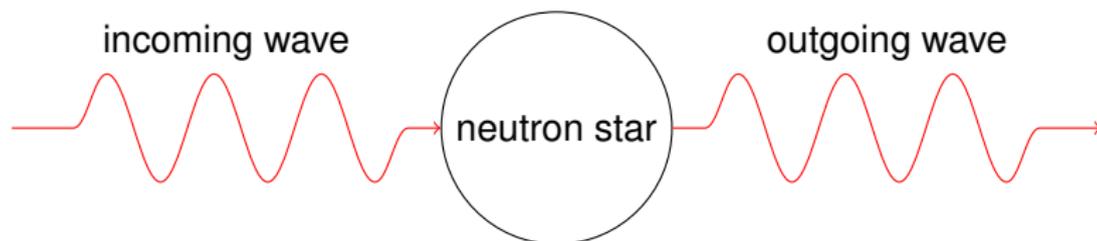
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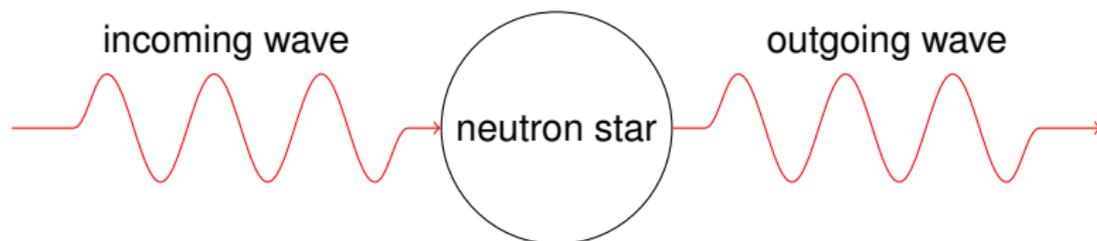
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$$\frac{d^2 X}{dr_*^2} + \left[\left(1 - \frac{2M}{r}\right) \frac{\ell(\ell + 1) - \frac{6M}{r}}{r^2} + \omega^2 \right] X = S$$

- How to construct solutions corresponding to external field and response?

Relativistic case: generic problems

- Analog to Hermitian operator \mathcal{D} not available in the relativistic case
- Substitute neutron star by point particle reproducing large-scale field

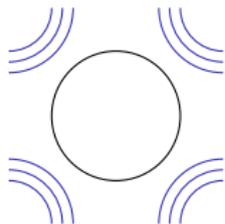


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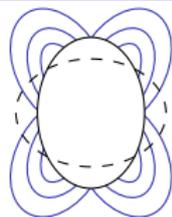
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- How to construct solutions corresponding to external field and response?

Identification of external field and response



external quadrupolar field



quadrupolar response

Newtonian:

$$r^{\ell+1}$$

adiabatic $\omega = 0$:

$$r^{\ell+1} {}_2F_1(\dots; 2M/r)$$

relativistic:

$$X_{\text{MST}}^{\ell}$$

$$r^{-\ell}$$

$$r^{-\ell} {}_2F_1(\dots; 2M/r)$$

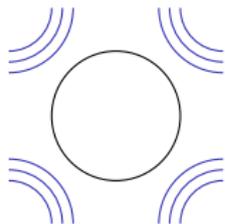
$$X_{\text{MST}}^{-\ell-1}$$

where [Mano, Suzuki, Takasugi, PTP **96** (1996) 549]

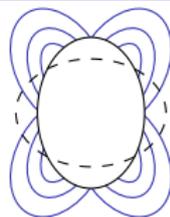
$$X_{\text{MST}}^{\ell} = e^{-i\omega r} (\omega r)^{\nu} \left(1 - \frac{2M}{r}\right)^{-i2M\omega} \sum_{n=-\infty}^{\infty} \dots \times \left[\frac{r}{2M}\right]^n {}_2F_1(\dots; 2M/r)$$

Renormalized angular momentum, transcendental number: $\nu = \nu(\ell, M\omega)$

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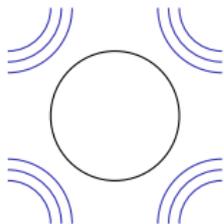
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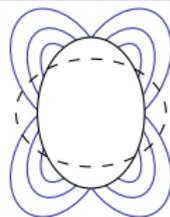
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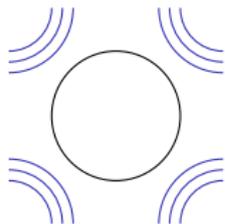
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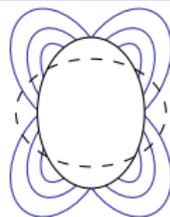
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Identification of external field and response by
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Relativistic response

Chakrabarti, Delsate, Steinhoff, arXiv:1304.2228

- Numerical neutron star perturbation matched to

$$X = A_1 X_{\text{MST}}^\ell + A_2 X_{\text{MST}}^{-\ell-1}$$

- X_{MST}^ℓ , $X_{\text{MST}}^{-\ell-1}$ related to effective point-particle source via
variation of parameters

- Point-particle source requires regularization (here: Riesz-kernel)
- Regularization introduces dependence on scale μ_0

- Fit for the response:

$$F(\omega) \approx \sum_n \frac{I_n^2}{\omega_n^2 - \omega^2}$$

- Just like in Newtonian case!
- Relative error less than 2%
- Relativistic overlap integrals: I_n
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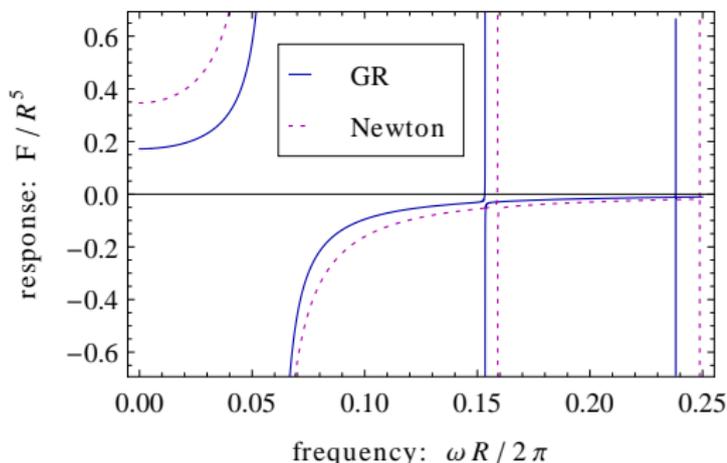
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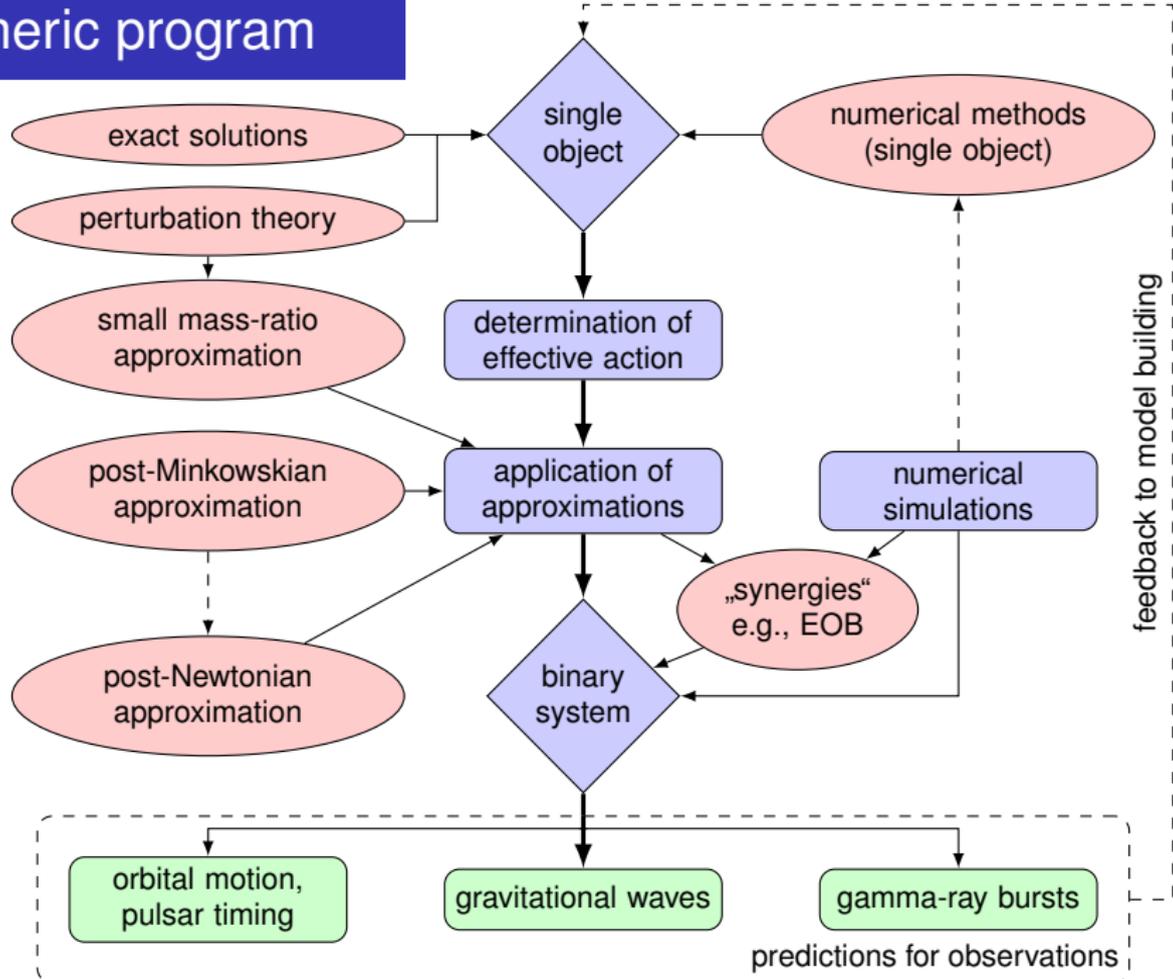
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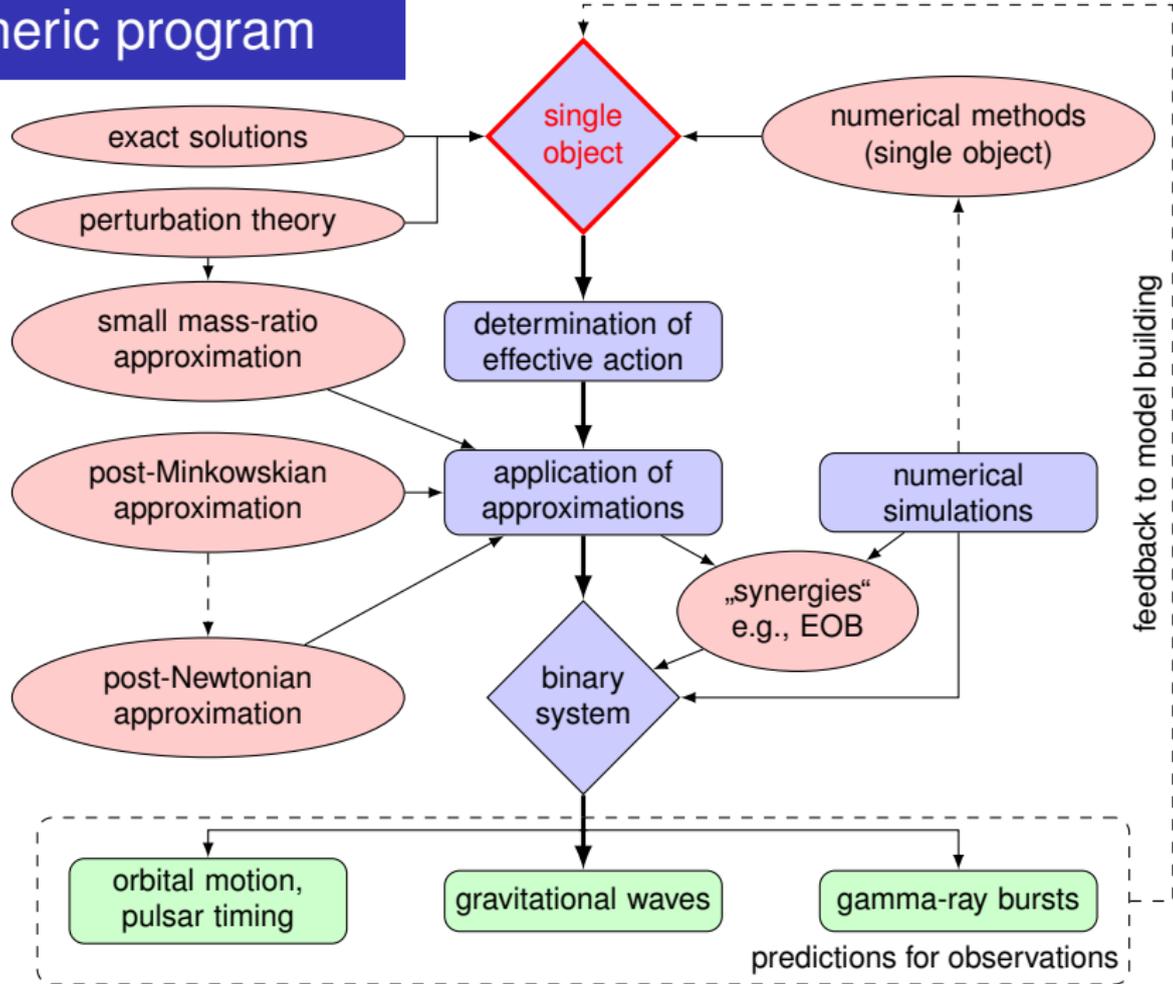
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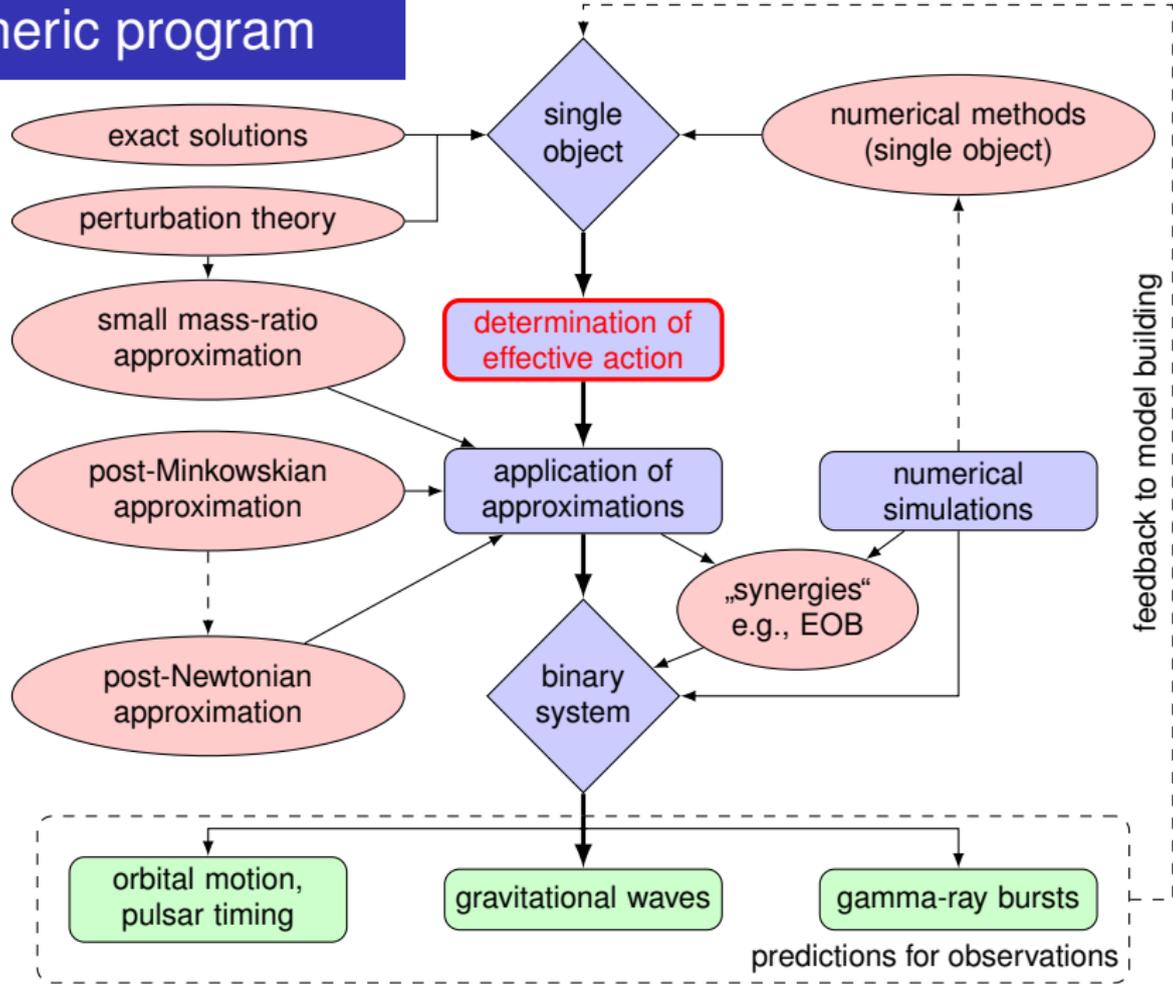
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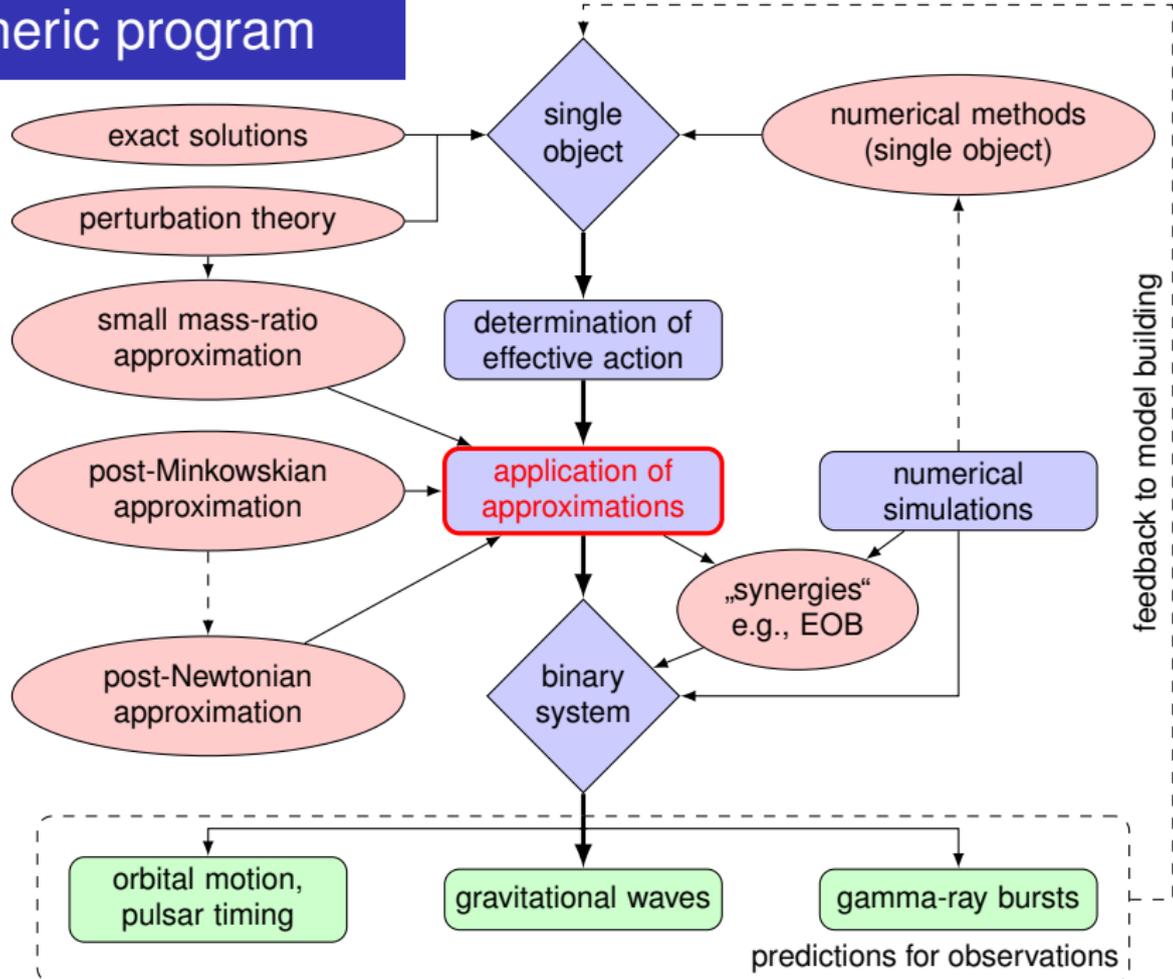
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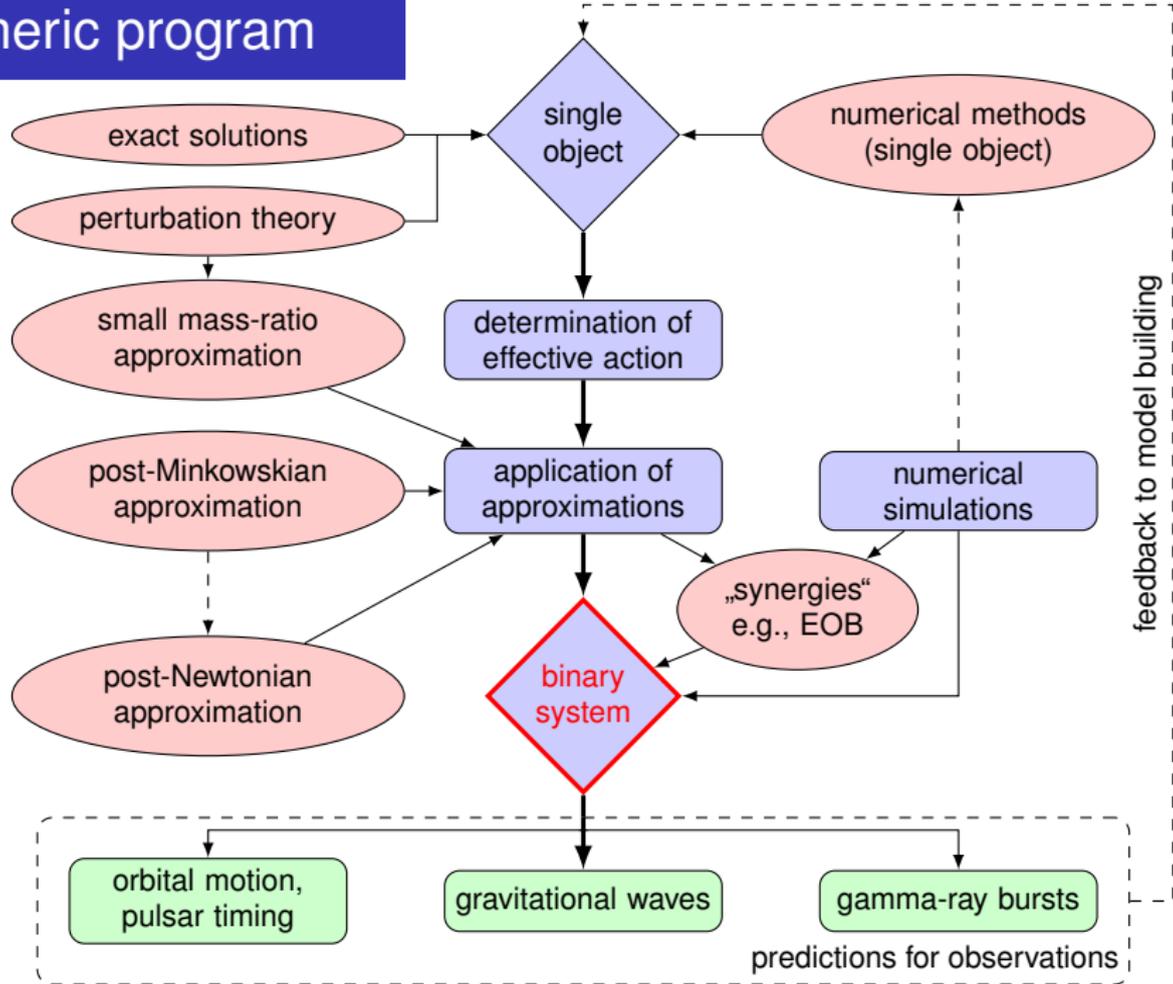
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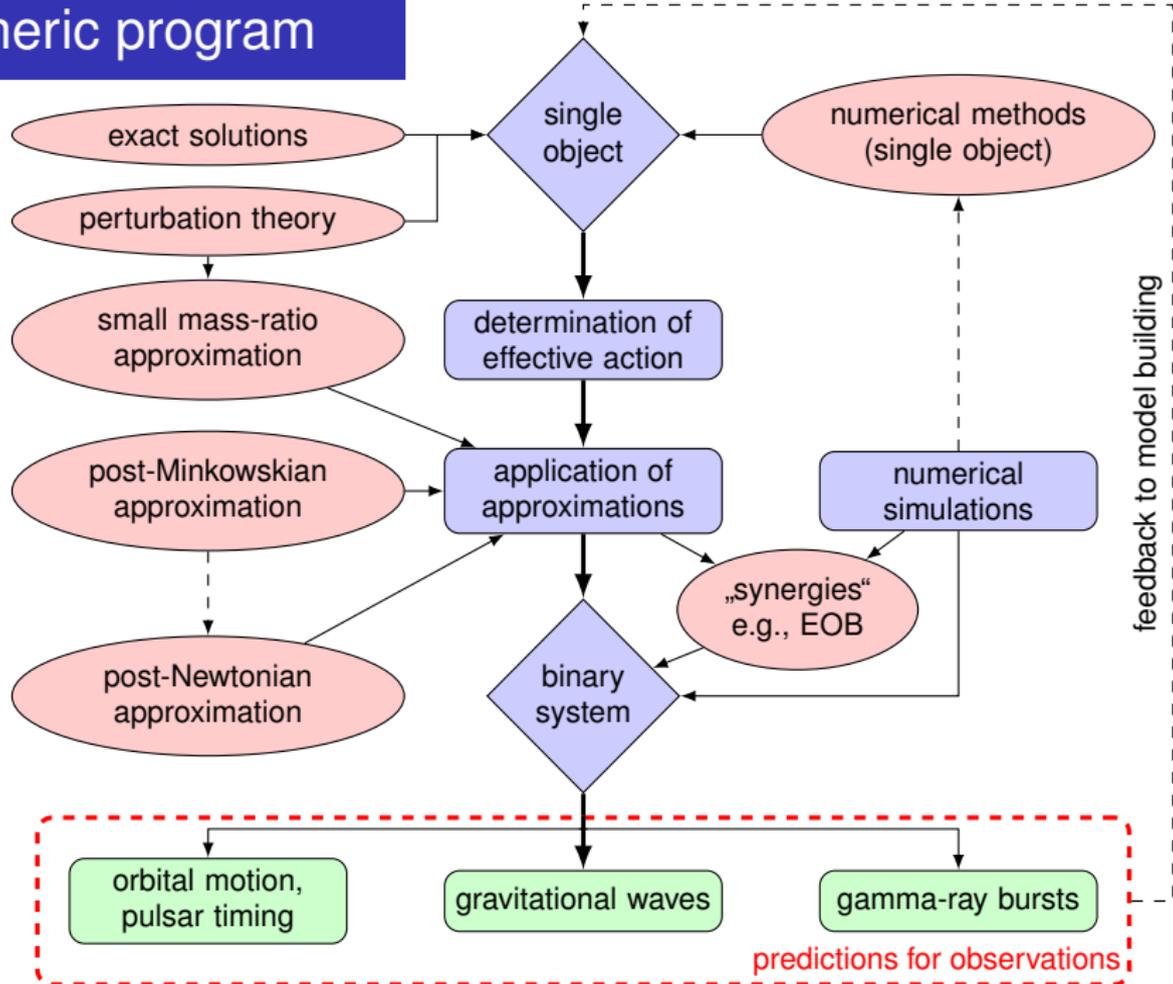
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Conclusions and outlook

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- We defined the relativistic quadrupolar response for linear perturbation of nonrotating barotropic stars
- Response is completely analogous to Newtonian case
- We defined **relativistic overlap integrals**
- Important step towards **gravito-spectroscopy** using gravitational waves

Outlook:

- More realistic neutron star models:
rotation, crust, . . . (also for Newtonian case)
- Connection to gamma-ray bursts:
shattering of crust, instabilities of modes
- Dimensional regularization
- Other multipoles
based on action in [Goldberger, Ross, PRD 81 (2010) 124015]
- 2nd Love number of **rotating** black holes

Thank you for your attention

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