

Analytic models for compact binaries

spin and **dynamic tides**

Jan Steinhoff

in collaboration with Tanja Hinderer, Alessandra Buonanno, and Andrea Taracchini

Steinhoff et al, arXiv:1608.01907

Hinderer et al, PRL **116** (2016) 181101



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Results for post-Newtonian approximation with spin

conservative part of the motion of the binary; see talk by Michele Levi on Monday

post-Newtonian (PN) approximation: expansion around $\frac{1}{c} \rightarrow 0$ (Newton)

order	c^0 N	c^{-1}	c^{-2} 1PN	c^{-3}	c^{-4} 2PN	c^{-5}	c^{-6} 3PN	c^{-7}	c^{-8} 4PN
non spin	✓		✓		✓		✓		✓
spin-orbit				✓		✓		✓	
S_1^2					✓		✓		✓
$S_1 S_2$					✓		✓		✓
Spin ³								✓ (✓)	
Spin ⁴									✓ (✓)
⋮									⋮
	✓ known		(✓) partial						✓ derived last year

Work by many people (“just” for the spin sector): Barker, Blanchet, Bohé, Buonanno, O’Connell, Damour, D’Eath, Faye, Hartle, Hartung, Hergt, Jaranowski, Marsat, Levi, Ohashi, Owen, Perrodin, Poisson, Porter, Porto, Rothstein, Schäfer, Steinhoff, Tagoshi, Thorne, Tulczyjew, Vaidya

Zones, separation of scales, and effective theory

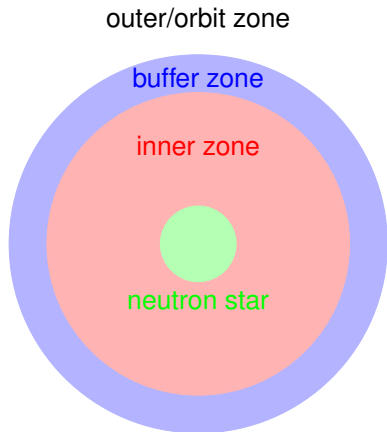
EFT program in classical gravity: Goldberger, Rothstein, PRD **73** (2006) 104029; ...

various zones → separation of scales

scales continue down the star:
→ fluid, nucleons, quarks, ?

The physics at “smaller” scales admits
an Effective Field Theory (EFT) description!

Here: Effective theory for dynamical tides
→ dynamical, time-dependent response
(of the inner zone to perturbations from the outer zone)
→ harmonic oscillator effective theory for multipoles



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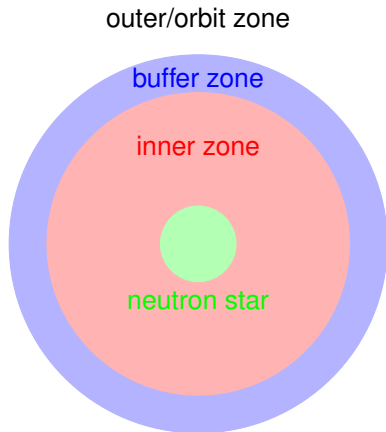
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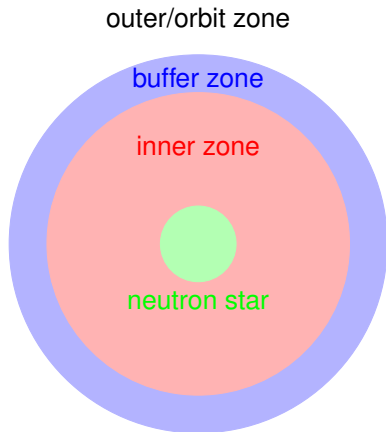
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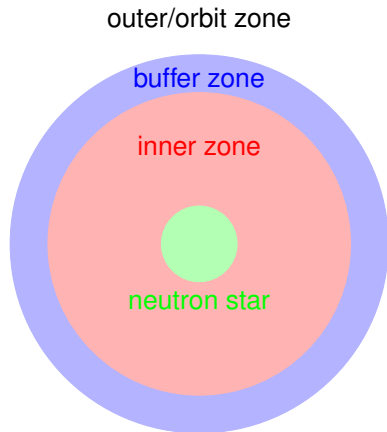
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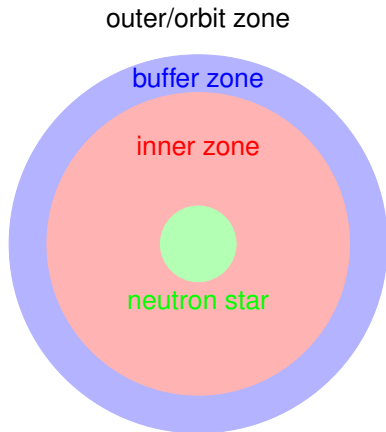
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Dynamical tides in general relativity

Their description through an effective action [JS, Hinderer, Taracchini, Buonanno, in preparation]

Relativistic effective Lagrangian for dynamical tides: $Q_{\mu\nu} u^\nu = 0$

$$L_Q = \frac{z}{4\lambda\omega_f^2} \left[\frac{1}{z^2} \frac{DQ_{\mu\nu}}{d\sigma} \frac{DQ^{\mu\nu}}{d\sigma} - \omega_f^2 Q_{\mu\nu} Q^{\mu\nu} \right] - \frac{z}{2} E_{\mu\nu} Q^{\mu\nu} + \frac{z}{4} K E_{\mu\nu} E^{\mu\nu} + \dots$$
$$u^\mu = \frac{Dx^\mu}{d\sigma}, \quad z = \sqrt{-u^\mu u_\mu} \quad (\text{is the redshift for } \sigma = t)$$

- Newtonian case: [Flanagan, Hinderer, PRD **77** (2008) 021502]
- λ is the tidal deformability (Love number)
- identify ω_f with real part of quasi-normal-mode frequency
- K linked to (almost) completeness of modes: $K \approx 0$

ω_f and K are not fixed by a matching, but by physical intuition!

a prescription for the dynamical response is in Chakrabarti, Delsate, JS, arXiv:1304.2228

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- redshift effect
- gravitomagnetism
 - frame dragging effect
 - ~ Zeeman effect

Both effectively shift
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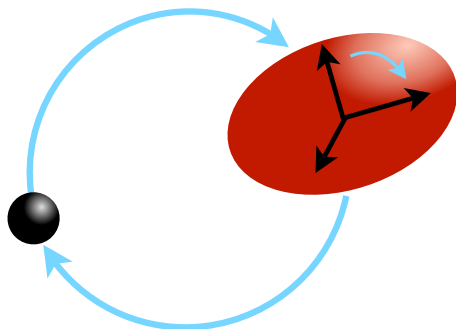
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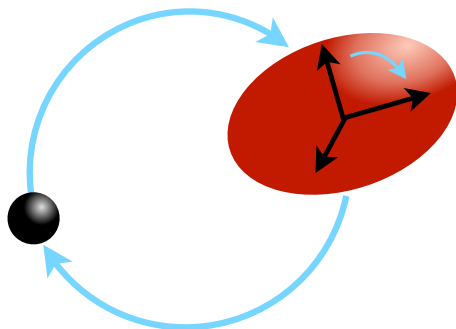
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Computing the post-Newtonian (PN) corrections

Frame dragging interaction

tidal spin: $S_Q^{ij} = 4Q^{k[i}P^{j]k}$
generates infinitesimal rotations
→ frame dragging

substitute $S^{ij} \rightarrow S_Q^{ij}$ in known potentials! → lazy

The tidal driving force

tidal: $-\frac{1}{2}E_{\mu\nu}Q^{\mu\nu}$ vs. spin induced: $\frac{C_{ES^2}}{2m}E_{\mu\nu}S^\mu S^\nu$

again substitute: $C_{ES^2}S^i S^j \rightarrow -mQ^{ij}$ in S^2 known potentials

super lazy!!! agrees with Vines, Flanagan, PD **88** (2013) 024046

Harder: implementation into effective-one-body, analyze various models, ...

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Conclusions

All you need is λ ! ?



Almost, need more coefficients
linked to dynamical tides!

$\lambda, \omega_f, K, \dots$

Dynamical tides become important
close to resonance with ω_f

Increase tidal effect by $\sim 30\%$!

Dynamical tides are important for accurate waveform models

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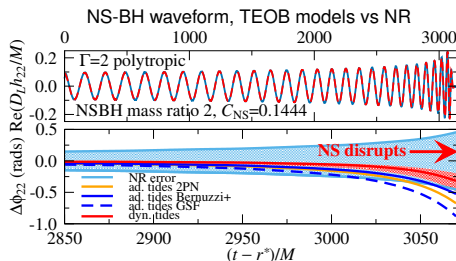
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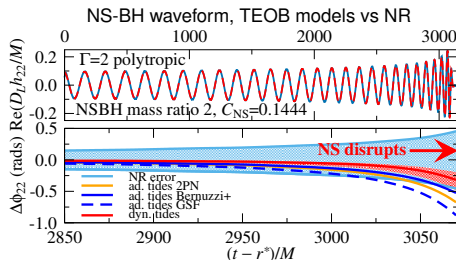
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