

Tidal forces and mode resonances in compact binaries

Sayan Chakrabarti¹ T rence Delsate² Jan Steinhoff³



¹Indian Institute of Technology
Guwahati, India



²UMons
Mons, Belgium



T CNICO
LISBOA

³Instituto Superior T cnico
Lisbon, Portugal

Seminar of the gravity group, Lisbon, January 8th, 2013

- 1 Motivation
 - Strong gravity
 - Neutron star
 - Tidal forces and resonances
- 2 Effective action and strong field gravity
 - Effective theory for compact objects
 - Philosophy
 - Example: adiabatic tidal deformation
- 3 Dynamic tidal interactions in Newtonian gravity
 - Tidal forces in Newtonian gravity
 - Effective theory point of view on tidal interactions
 - Convenient concept: response function
 - Analogy with electronics
- 4 Relativistic case
 - Identification of external field and response
 - Relativistic response
- 5 Outlook
 - Outlook

Strong gravity and compact objects

General relativity is expected to break down for strong field strengths.

Currently our best chance to test strong gravity:

compact astrophysical objects!

Black holes: clean (vacuum solutions); but observations difficult

Neutron stars (NS):

- strong gravity in NS interior
- alternative theories can predict drastically different NS structure: oscillation modes, mass-radius relation, spontaneous scalarization
but partly unknown matter properties . . .

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Neutron star picture by D. Page

www.astroscu.unam.mx/neutrones/

„Lab“ for various areas in physics

- magnetic field, plasma
- crust (solid state)
- superfluidity
- superconductivity
- unknown matter in core
condensate of quarks, hyperons,
kaons, pions, ... ?
accumulation of dark matter ?

pics/neutronstar

Related objects:

- pulsars
- magnetars
- quark stars

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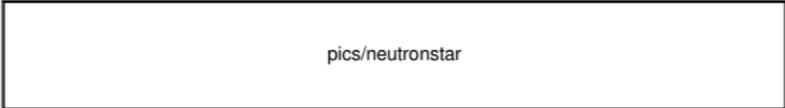
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Tidal forces in binaries and resonances

tidal forces in inspiraling binaries \longleftrightarrow oscillation modes of neutron stars
 \Rightarrow resonances!

pics/doublepulsar

binary neutron stars

Swift/BAT, nasa.gov

resonances probably
relevant for short
gamma-ray bursts

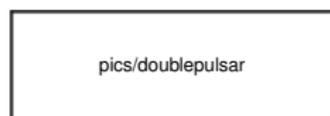
[Tsang et.al., PRL 108 (2012) 011102]

mode spectrum from
gravitational waves:
gravito-spectroscopy

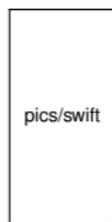
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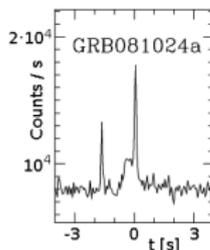
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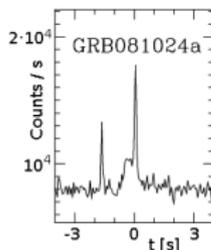
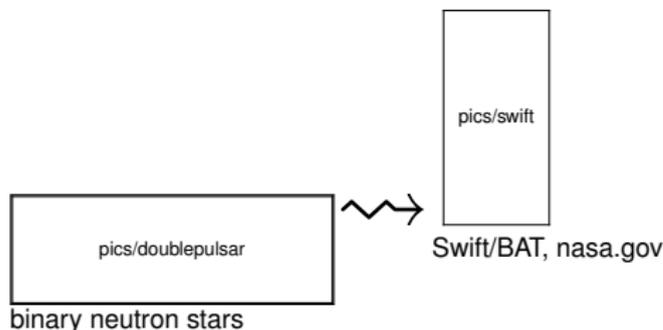
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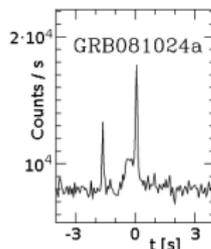
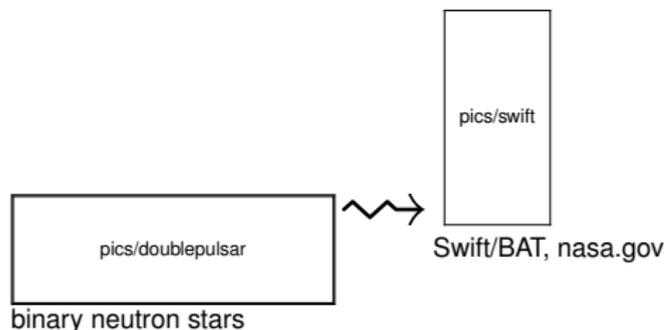
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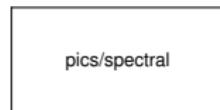
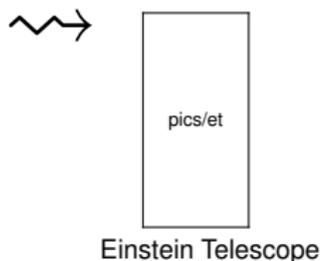
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Effective theory for compact objects

See: [Goldberger, Rothstein, PRD **73** (2006) 104029]

[Goldberger, Ross, PRD **81** (2010) 124015]

Effective theory: important idea from statistical physics, quantum field theory

Starting point: An action for a “full theory” S_{full}

“Integrate out” short scale (UV) part of the metric g_{UV} :

$$\exp\left(\frac{i}{\hbar} S_{\text{eff}}[g_{IR}, \text{matter}]\right) = \int Dg_{UV} \exp\left(\frac{i}{\hbar} S_{\text{full}}[g_{IR} + g_{UV}, \text{matter}]\right)$$

We are only interested in the **classical part** of the path integral.

On large scales (IR), a mass distribution looks like a **point-particle**:

$$S_{\text{eff,matter}}[g_{IR}, \text{matter}] = \int d\tau \left[-m - \frac{1}{2} E_{ab} Q^{ab} + \dots \right]$$

$$E_{ab} \sim R_{acbd} U^c U^d, \quad Q^{ab} \sim \text{mass quadrupole}$$

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Philosophy of the approach

Starting point: single object, e.g., neutron star

Idea

Multipoles describe compact object on macroscopic scale

\longleftrightarrow

state variables (ρ, V, T)	\longleftrightarrow	multipoles (m, S, Q)
equations of state	\longleftrightarrow	effective action
correlation	\longleftrightarrow	response

Approximations for binary system using effective theory:

- Effective point-particle action with dynamical multipoles
- Response functions (propagators) for multipoles

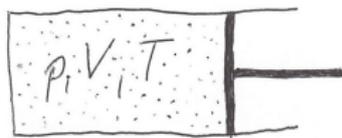
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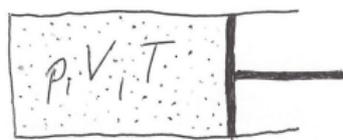
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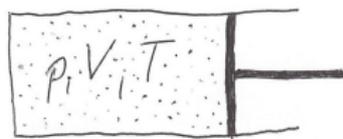
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Example: **adiabatic** tidal deformation

for NS, e.g. Hinderer & Flanagan (2008); Damour, Nagar (2009); Binnington, Poisson (2009)

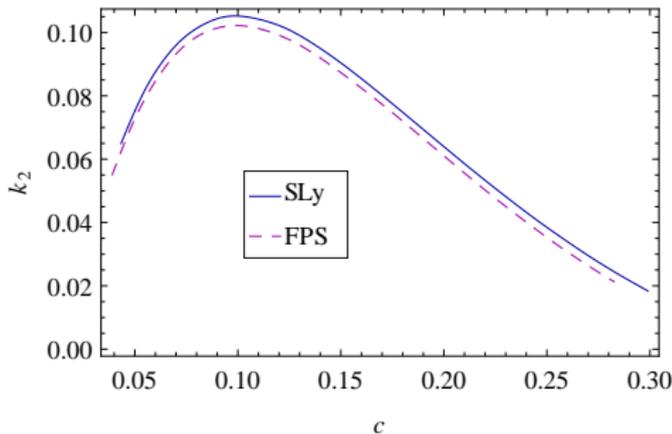
- **Static**, linear NS perturbation:

$$-Q = \mu_2 E$$

- Tidal force E (curvature)
- Dim.-less 2nd Love number k_2 :

$$k_2 = \frac{3}{2} \frac{\mu_2}{R^5}$$

- Measure for grav. polarizability
- Compactness $c = \frac{Gm}{R}$



see Damour, Nagar, PRD **80** (2009) 084035

- k_2, μ_2 encode **strong gravity** effects in the interior
- beyond the adiabatic case:

[Bini, Damour, Faye, PRD **85** (2012) 124034]

[Maselli, Gualtieri, Pannarale, Ferrari, PRD **86** (2012) 044032]

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Tidal forces in Newtonian gravity

Simple case: **linear perturbation of nonrotating barotropic stars**

temperature-independent equation of state

see e.g. [Press, Teukolsky, ApJ **213** (1977) 183]

Displacement $\vec{\xi}$:= perturbed – unperturbed location of fluid elements

$$\ddot{\vec{\xi}} + \mathcal{D}\vec{\xi} = (\text{external forces})$$

$$\mathcal{D}\vec{\xi} := -\vec{\nabla} \left\{ \left[\frac{c_s^2}{\rho_0} + 4\pi G\Delta^{-1} \right] \vec{\nabla} \cdot (\rho_0 \vec{\xi}) \right\}$$

ρ_0 : unperturbed mass density

c_s : speed of sound

G : Newton constant

The operator \mathcal{D} :

- differential $\vec{\nabla}$ and integral Δ^{-1} operators
- linear, nonlocal, spherical symmetric
- **Hermitian w.r.t. compact measure** $dm_0 := \rho_0 d^3x$

[Chandrasekhar, ApJ **139** (1964) 664–674]

⇒ Eigenfunctions of \mathcal{D} are the normal oscillation modes of the star with discrete, real eigenvalues $\omega_{n\ell}^2 \rightarrow$ oscillation frequencies

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Effective theory point of view on tidal interactions

Chakrabarti, Delsate, Steinhoff, PRD **88** (2013) 084038

Full theory: variational fluid dynamics, linear perturbation

Integrate out small scales, **mode decomposition**:

$$\vec{\xi} = \sum_{n\ell m} A_{n\ell m}(t) \vec{\xi}_{n\ell m}^{\text{NM}}(\vec{x})$$

Result for effective action, **quadrupolar truncation** $\ell = 2$:

$$S_{\text{eff,matter}} = \int dt \left[\frac{1}{2} m \dot{\vec{z}}^2 - m\Phi + \frac{1}{2} \sum_{n,m} \left[|\dot{A}_{n2m}|^2 - \omega_n^2 |A_{n2m}|^2 - I_n A_{n2m} E_{2m} \right] + \dots \right]$$

$E_{2m} \sim E_{ab} = \partial_a \partial_b \Phi, \quad n = \text{mode number}$

Overlap integrals I_n : coupling constants

Gravito-spectrum: ω_n and I_n

Try this with a nonlinear extension of Newtonian gravity?

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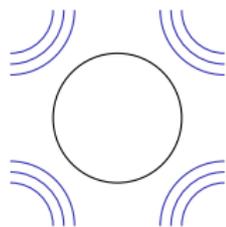
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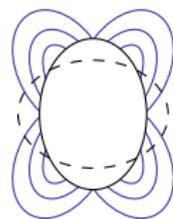
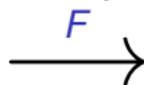
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external quadrupolar field

$$\Phi \sim r^\ell$$

linear response



quadrupolar response

$$\Phi \sim r^{-\ell-1}$$

→ deformation →

quadrupolar
response:

poles ⇒ resonances!

$$F = \sum_n \frac{I_n^2}{\omega_n^2 - \omega^2}$$

ω_n : mode frequency

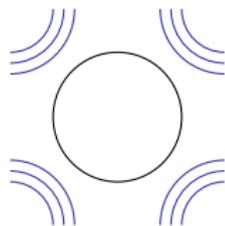
I_n : overlap integral

R : radius

Computation of I_n through fit of F ⇒ generalizes to relativistic case

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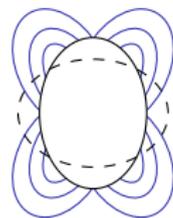
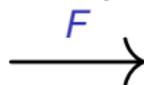
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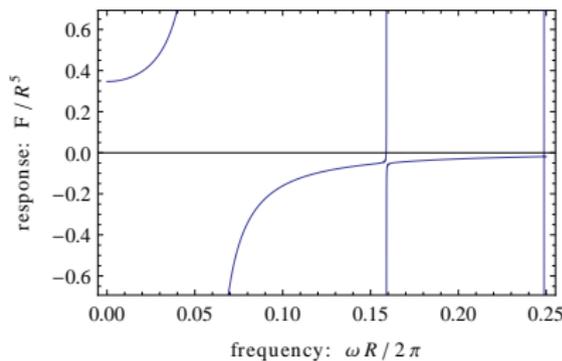
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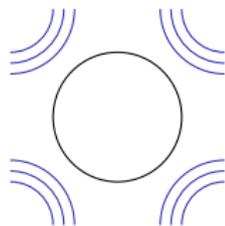
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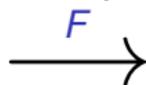
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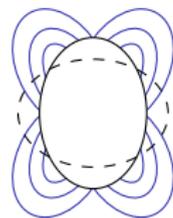
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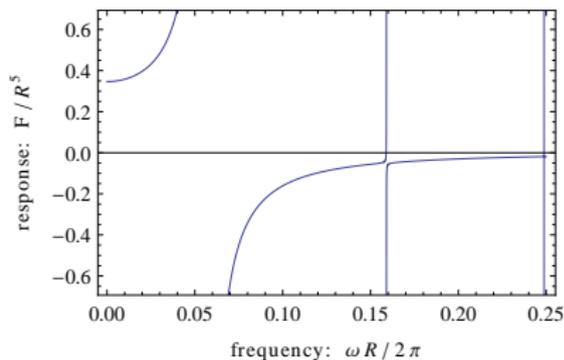


→ deformation →



quadrupolar response

$$\Phi \sim r^{-\ell-1}$$



quadrupolar response:

$$F = \sum_n \frac{I_n^2}{\omega_n^2 - \omega^2}$$

ω_n : mode frequency

I_n : overlap integral

R : radius

poles \Rightarrow resonances!

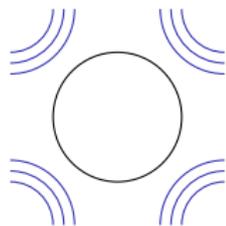


Tacoma Bridge

Computation of I_n through fit of $F \Rightarrow$ generalizes to relativistic case

Convenient concept: response function

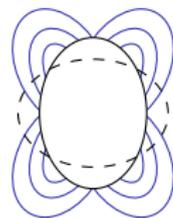
Chakrabarti, Delsate, Steinhoff, PRD **88** (2013) 084038



external quadrupolar field

$$\Phi \sim r^\ell$$

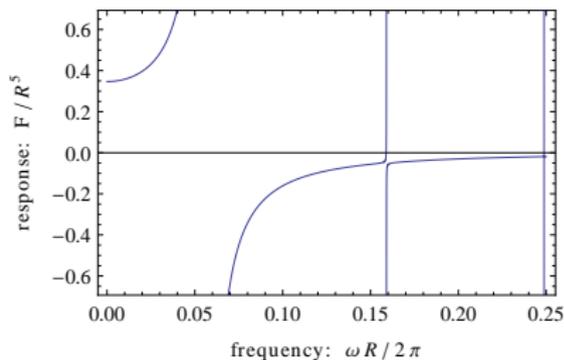
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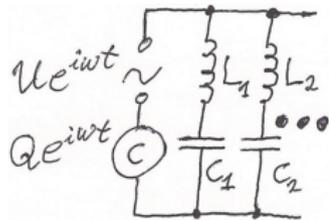
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Computation of I_n through fit of $F \Rightarrow$ generalizes to relativistic case

Analogy with electronics

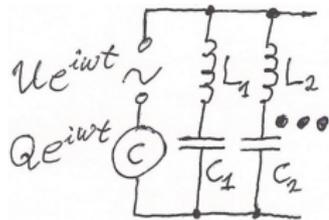


$$\begin{aligned}\frac{Q}{U} &= \frac{1}{i\omega Z} \\ &= \sum_n \frac{\frac{1}{L_n}}{\frac{1}{C_n L_n} - \omega^2}\end{aligned}$$

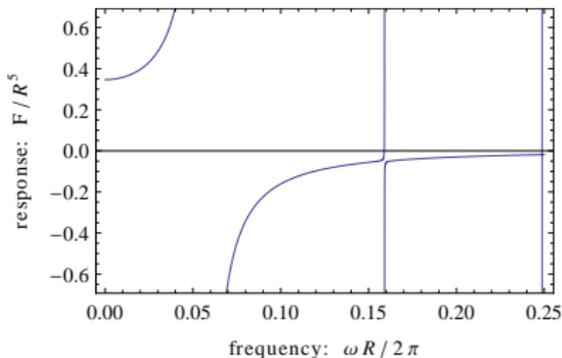
$$\begin{aligned}\frac{Q}{E} &=: F \\ &= \sum_n \frac{f_{nl}^2}{\omega_{nl}^2 - \omega^2}\end{aligned}$$

Q: quadrupole
E: external tidal field

Analogy with electronics



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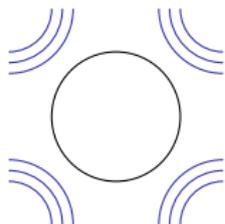


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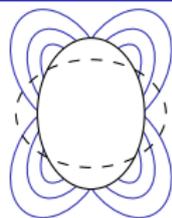
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Identification of external field and response



external quadrupolar field



quadrupolar response

Newtonian:

$$r^{\ell+1}$$

relativistic, adiabatic $\omega = 0$:

$$r^{\ell+1} {}_2F_1(\dots; 2m/r)$$

relativistic, generic ω :

$$X_{\text{MST}}^{\ell}$$

$$r^{-\ell}$$

$$r^{-\ell} {}_2F_1(\dots; 2m/r)$$

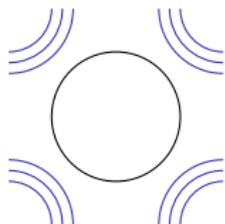
$$X_{\text{MST}}^{-\ell-1}$$

where [Mano, Suzuki, Takasugi, PTP 96 (1996) 549]

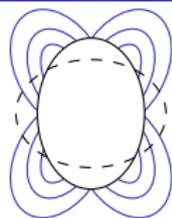
$$X_{\text{MST}}^{\ell} = e^{-i\omega r} (\omega r)^{\nu} \left(1 - \frac{2m}{r}\right)^{-i2m\omega} \sum_{n=-\infty}^{\infty} \dots \times \left[\frac{r}{2m}\right]^n {}_2F_1(\dots; 2m/r)$$

Renormalized angular momentum, transcendental number: $\nu = \nu(\ell, m\omega)$

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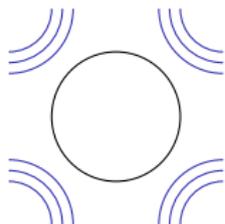
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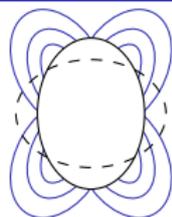
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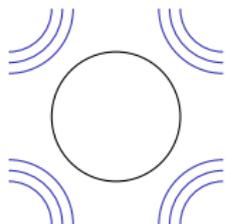
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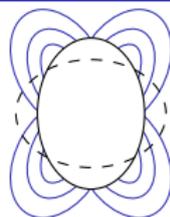
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Identification of external field and response by considering **generic** ℓ (analytic continuation)

- Numerical neutron star perturbation matched to

$$X = A_1 X_{\text{MST}}^\ell + A_2 X_{\text{MST}}^{-\ell-1} \quad (\text{homogeneous solution})$$

- $X_{\text{MST}}^\ell, X_{\text{MST}}^{-\ell-1}$ related to effective point-particle source via
variation of parameters (inhomogeneous solution)

- Point-particle source requires regularization (here: Riesz-kernel)

- Regularization introduces dependence on scale μ_0

- Fit for the response:

$$F(\omega) \approx \sum_n \frac{I_n^2}{\omega_n^2 - \omega^2}$$

- Just like in Newtonian case!
- Relative error less than 2%
- Relativistic overlap integrals: I_n
- Matching scale μ_0 is fitted, too

Relativistic response

Chakrabarti, Delsate, Steinhoff, arXiv:1304.2228

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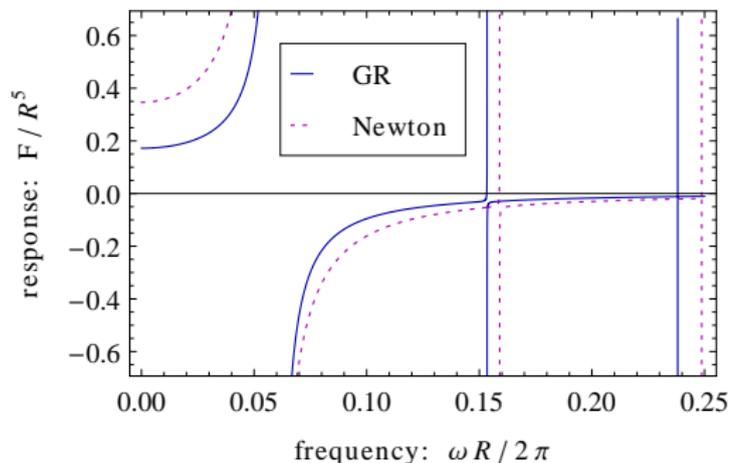
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- More realistic neutron star models: rotation, crust, ...
- Connection to gamma-ray bursts: shattering of crust

The search for universality among different NS models:

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